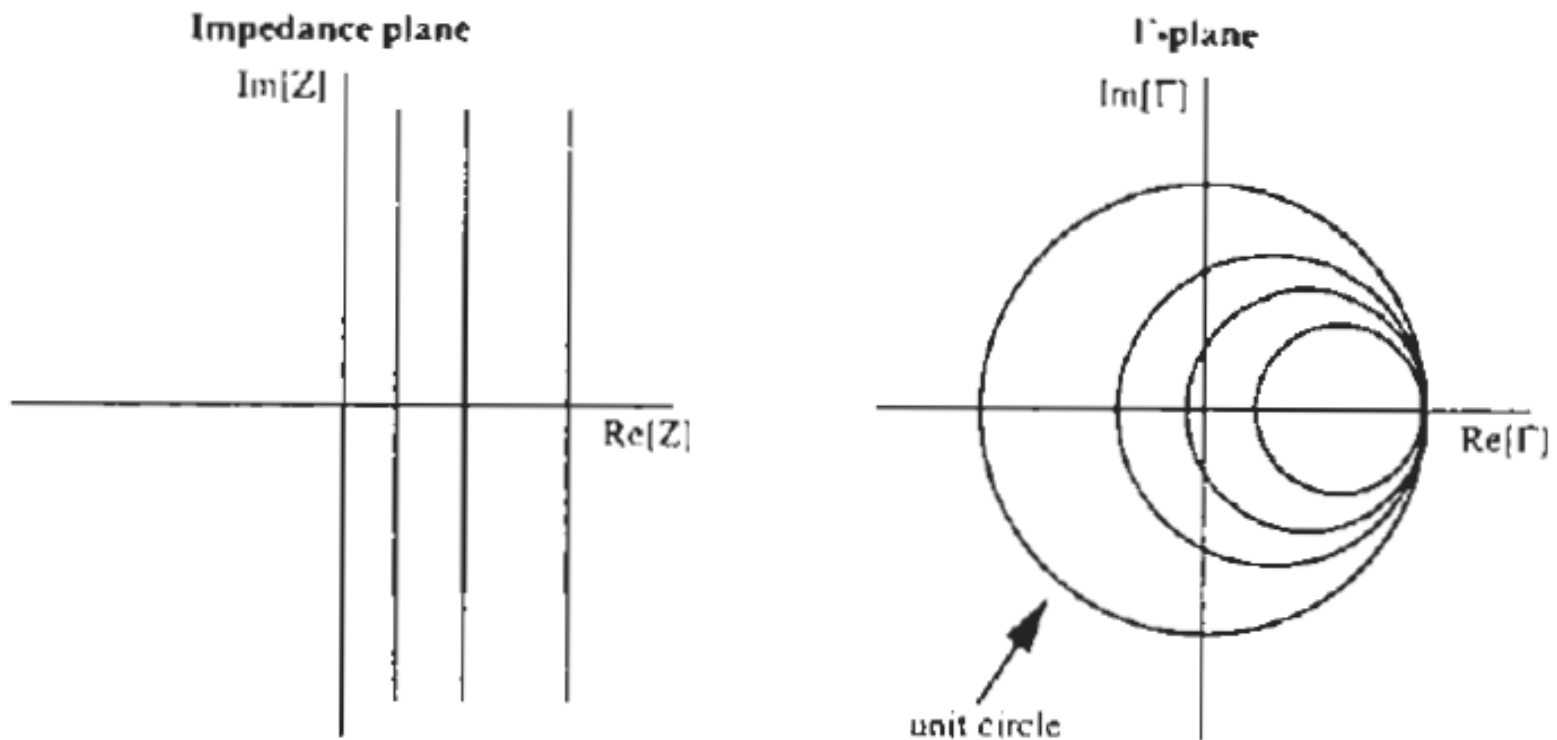


The Smith Chart and S-Parameters

Reflection Coefficient

$$\Gamma = \frac{z_l - z_o}{z_l + z_o} = \frac{\frac{z_l}{z_o} - 1}{\frac{z_l}{z_o} + 1} = \frac{z_{nl} - 1}{z_{nl} + 1}$$

Constant Resistance



Mapping of constant-resistance lines in Z -plane to circles in Γ -plane.

Proof:

assume: $Z_0 = R_0 + jX_0$, $Z_L = R_L + jX_L \rightarrow z_{nl} = r_l + jx_l = Z_L / Z_0$

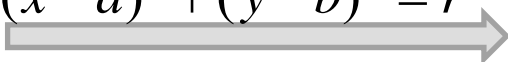
$$\Gamma = \frac{z_l - z_0}{z_l + z_0} = \frac{\frac{z_l}{z_0} - 1}{\frac{z_l}{z_0} + 1} = \frac{z_{nl} - 1}{z_{nl} + 1}$$

$$\left. \begin{aligned} z_{nl}\Gamma + \Gamma &= z_{nl} - 1 \\ z_{nl} &= \frac{1 + \Gamma}{1 - \Gamma} \end{aligned} \right\} \rightarrow r_l + jx_l = \frac{1 + \Gamma_r + j\Gamma_i}{1 - \Gamma_r - j\Gamma_i}$$

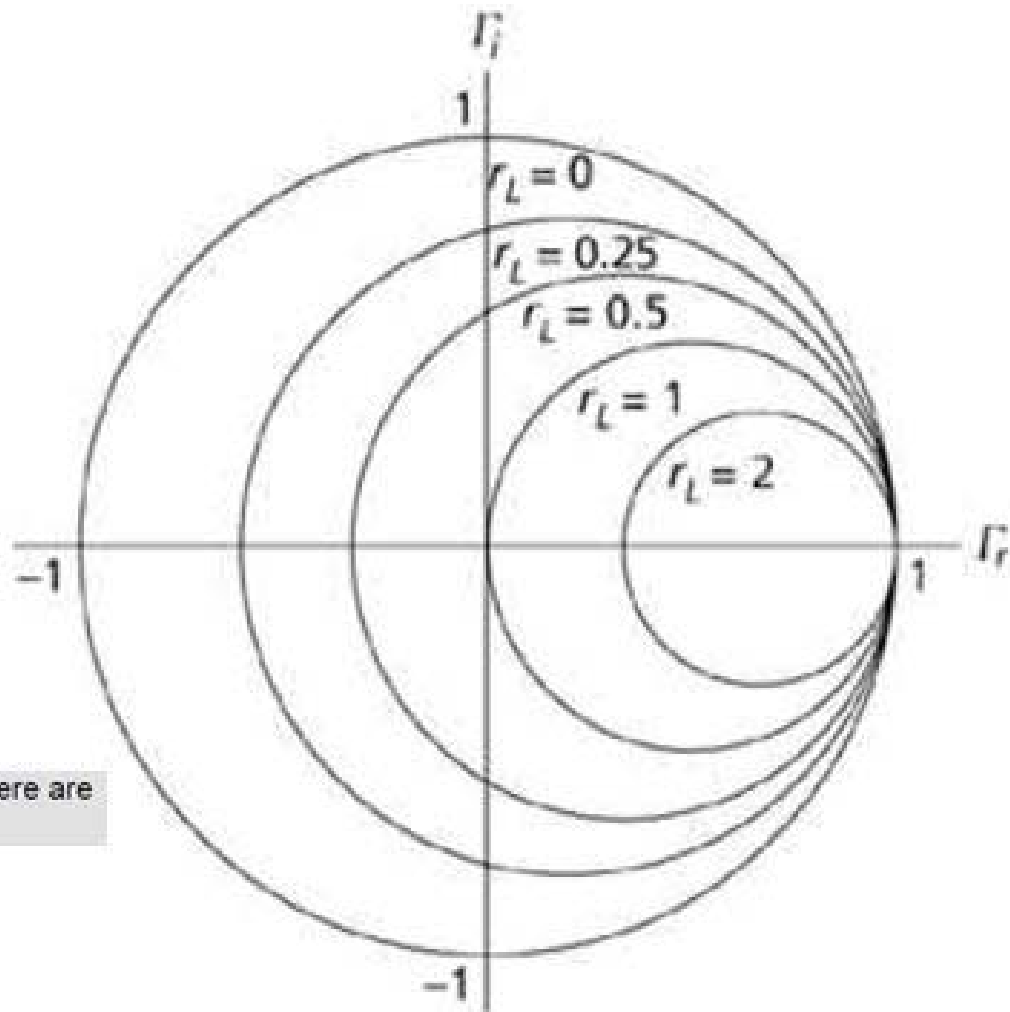
$$r_l + jx_l = \frac{1 + \Gamma_r + j\Gamma_i}{1 - \Gamma_r - j\Gamma_i}$$

$$r_l + jx_l = \frac{1 - \Gamma_r^2 - \Gamma_i^2}{(1 - \Gamma_r)^2 + \Gamma_i^2} + j \frac{2\Gamma_i}{(1 - \Gamma_r)^2 + \Gamma_i^2}$$

$$r_l = \frac{1 - \Gamma_r^2 - \Gamma_i^2}{\Gamma_r^2 - 2\Gamma_r + 1 + \Gamma_i^2}$$

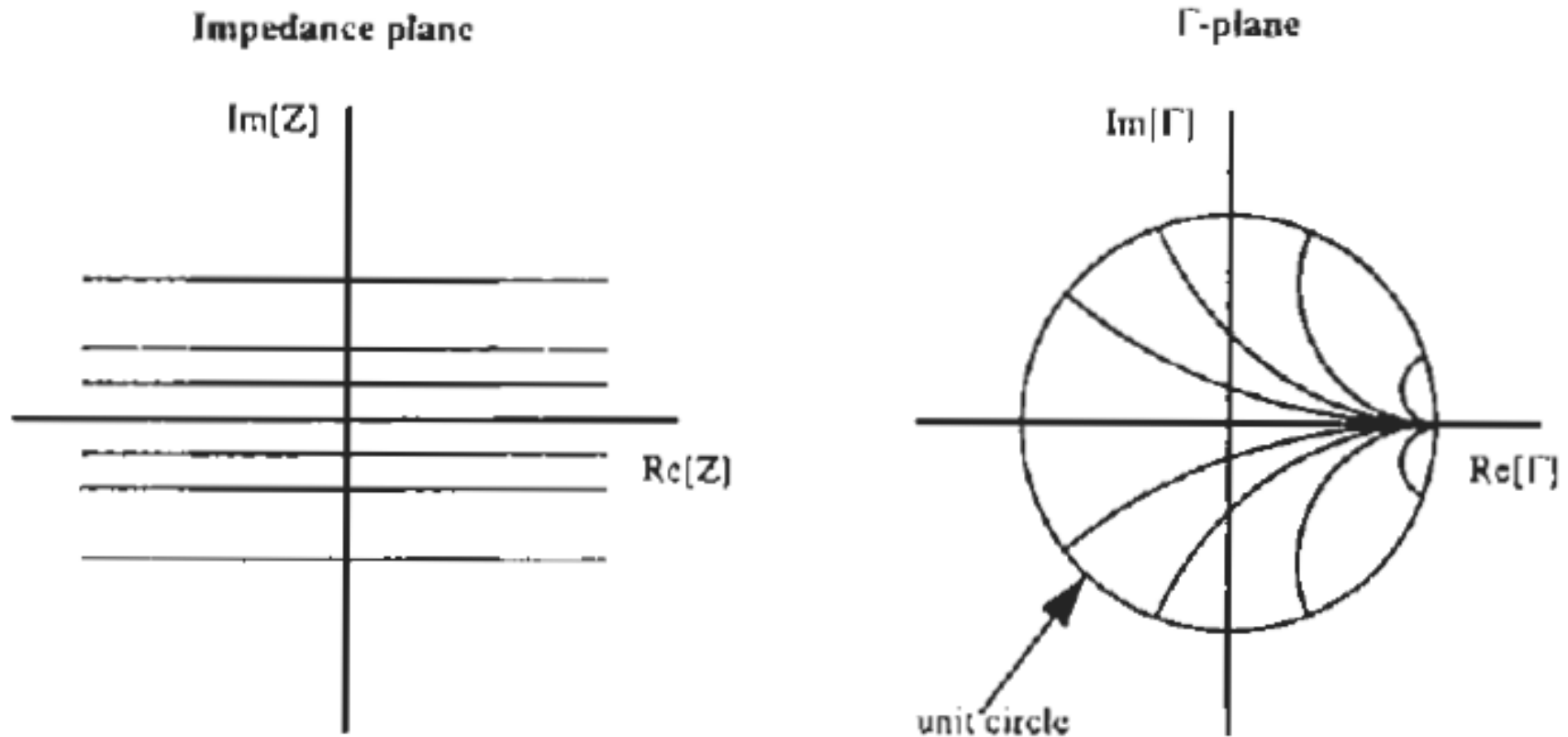
$$(x - a)^2 + (y - b)^2 = r^2$$


$$\left[\Gamma_r - \frac{r_l}{r_l + 1} \right]^2 + [\Gamma_i - 0]^2 = \left[\frac{1}{r_l + 1} \right]^2$$



Points of constant resistance form circles on the complex reflection-coefficient plane. Shown here are the circles for various values of load resistance.

Constant Reactance



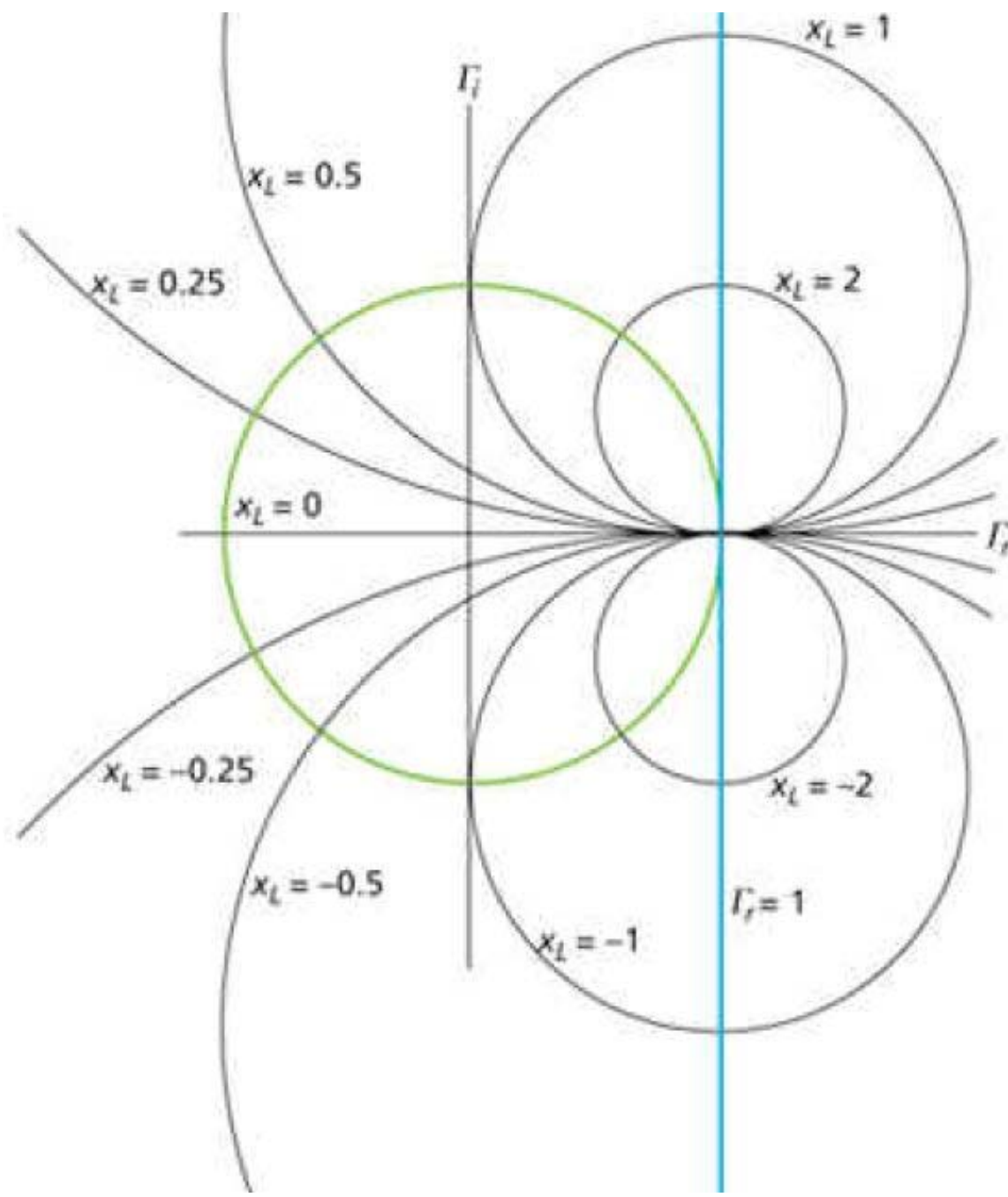
Mapping of constant-reactance lines in Z -plane to contours in Γ -plane.

$$r_l + jx_l = \frac{1 - \Gamma_r^2 - \Gamma_i^2}{(1 - \Gamma_r)^2 + \Gamma_i^2} + j \frac{2\Gamma_i}{(1 - \Gamma_r)^2 + \Gamma_i^2}$$

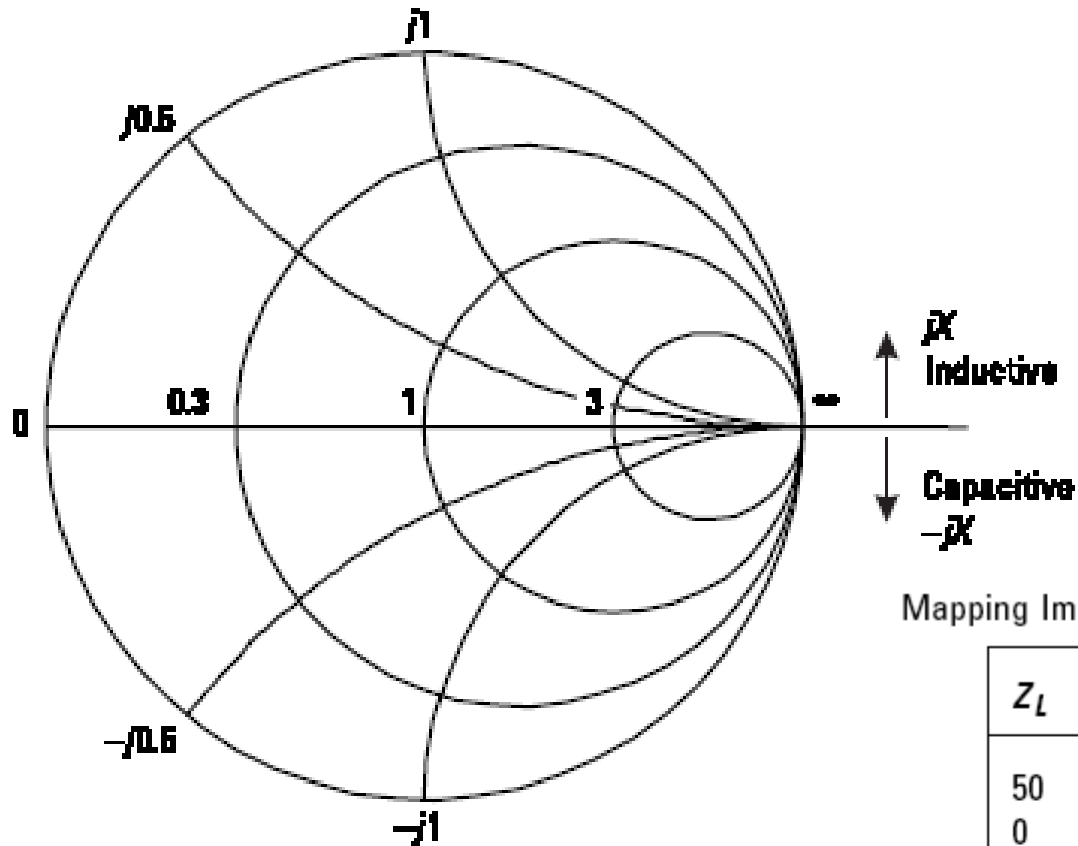
$$x_l = \frac{2\Gamma_i}{1 - 2\Gamma_r + \Gamma_r^2 + \Gamma_i^2}$$

$$(x - a)^2 + (y - b)^2 = r^2 \longrightarrow$$

$$(\Gamma_r - 1)^2 + \left(\Gamma_i - \frac{1}{x_l}\right)^2 = \left(\frac{1}{x_l}\right)^2$$

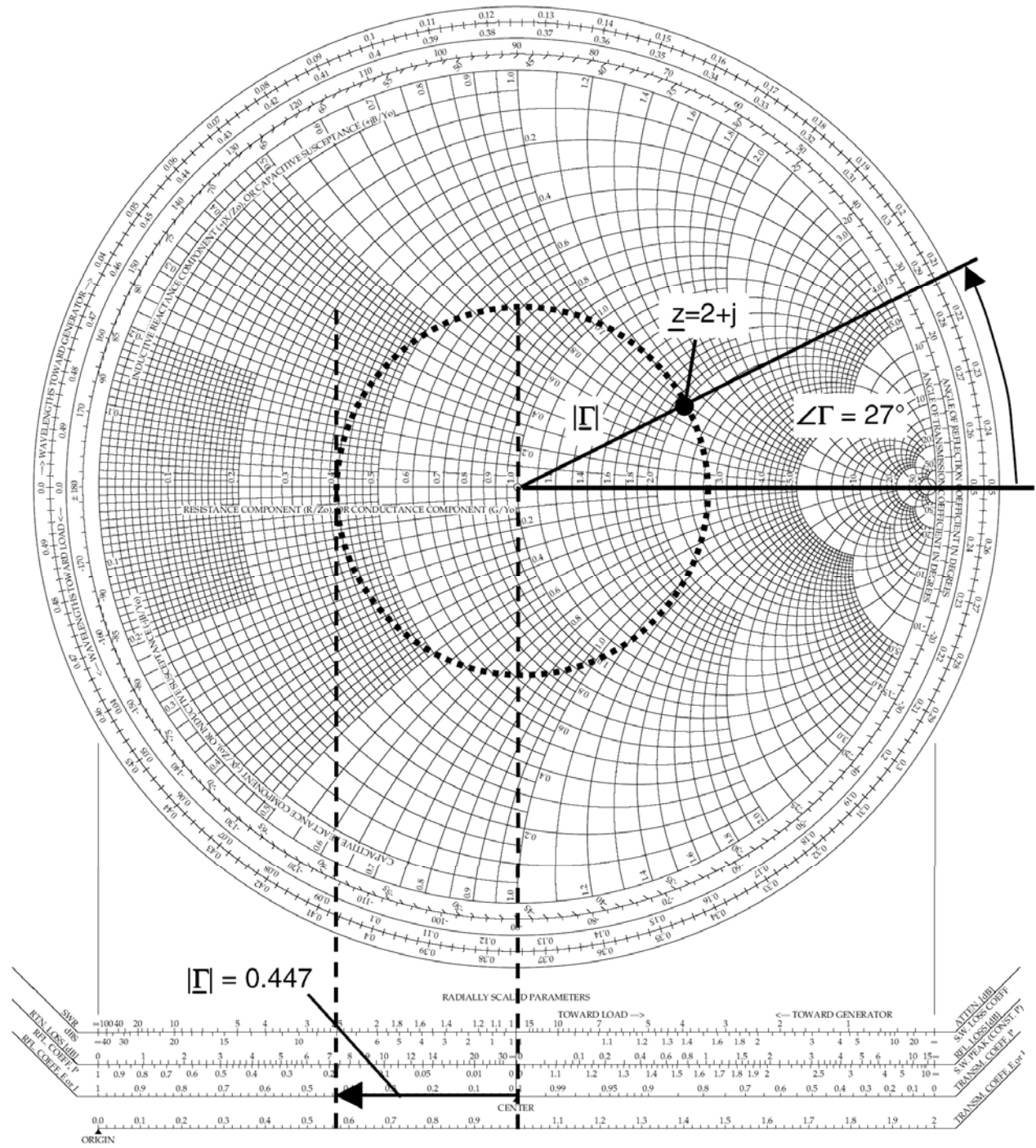


A Z Smith Chart



Mapping Impedances to Points on the Smith Chart

Z_L	Γ
50	0
0	-1
∞	1
100	0.333
25	-0.333
$j50$	$1 \angle 90^\circ$
jX	$1 \angle 2 \tan^{-1}(X/50)$
$50 - j141.46$	$0.8166 \angle -35.26^\circ$



For this simple example we can calculate

$$\underline{\Gamma} = \frac{\underline{Z} - Z_0}{\underline{Z} + Z_0} = \frac{\underline{z} - 1}{\underline{z} + 1} = \frac{2 + j - 1}{2 + j + 1} = \frac{1 + j}{3 + j}$$

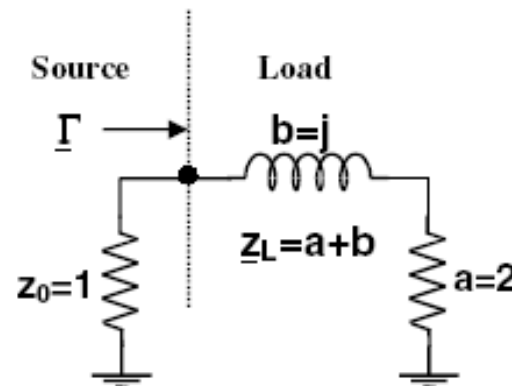
Separated into magnitude and phase, $\underline{\Gamma}$ is given by

$$|\underline{\Gamma}| = 0.447$$

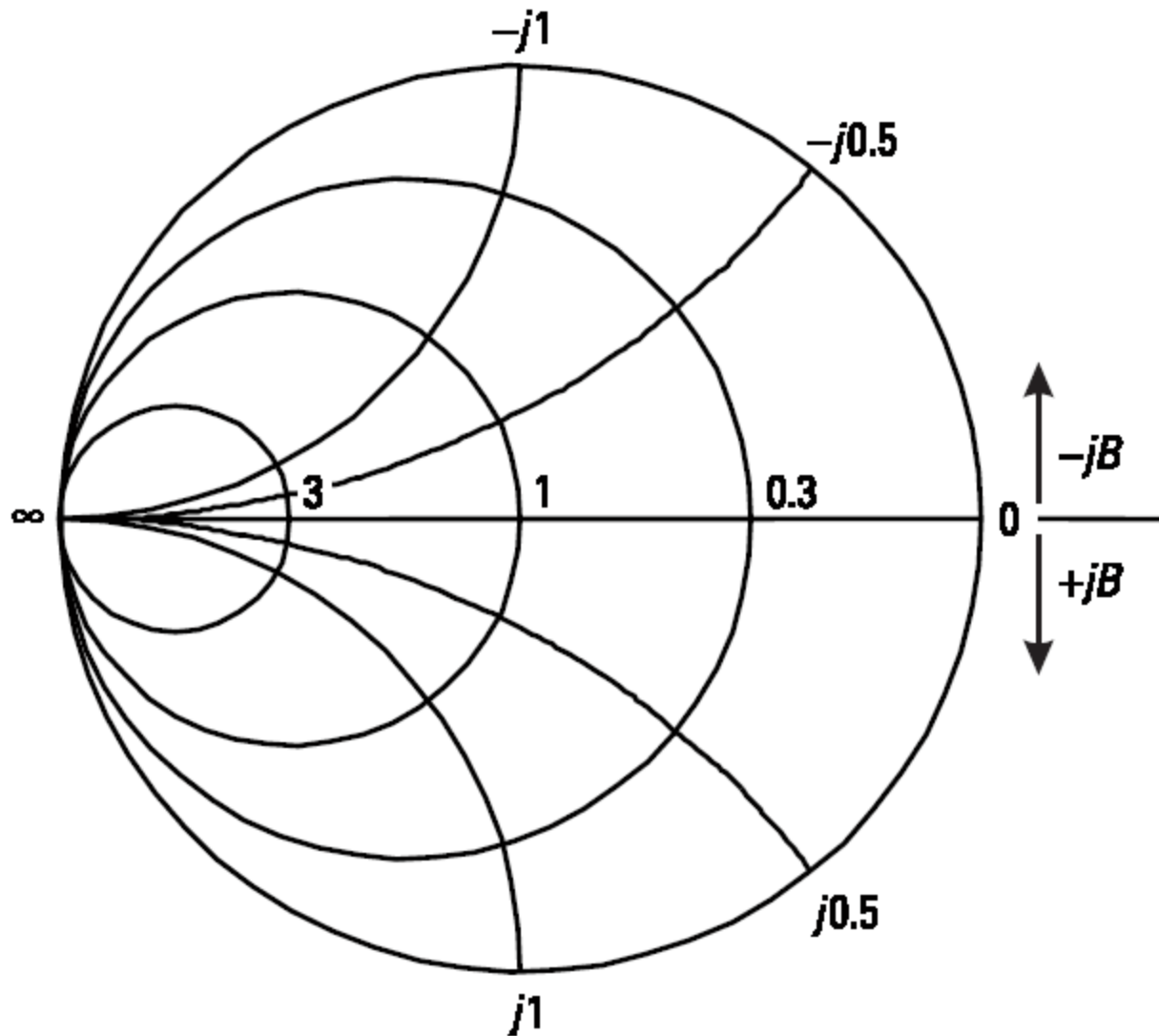
and

$$\angle \Gamma = \arctan(1) - \arctan\left(\frac{1}{3}\right) = 27^\circ$$

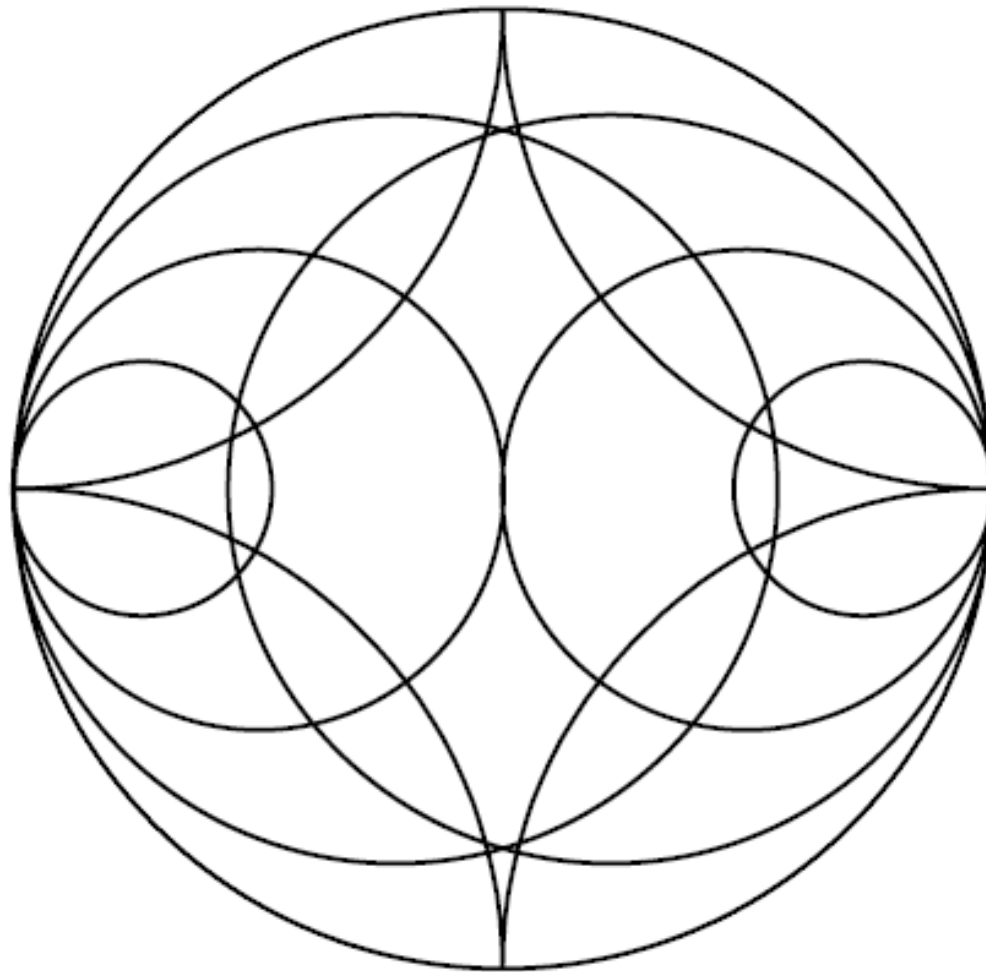
Consequently, the magnitude of the reflected power is proportional to $|\underline{\Gamma}|^2$. In our example, 20% of the power is reflected at the impedance transition, whereas the remaining 80% ($1 - |\underline{\Gamma}|^2$) is fed into the load.



An Admittance or Y Smith Chart

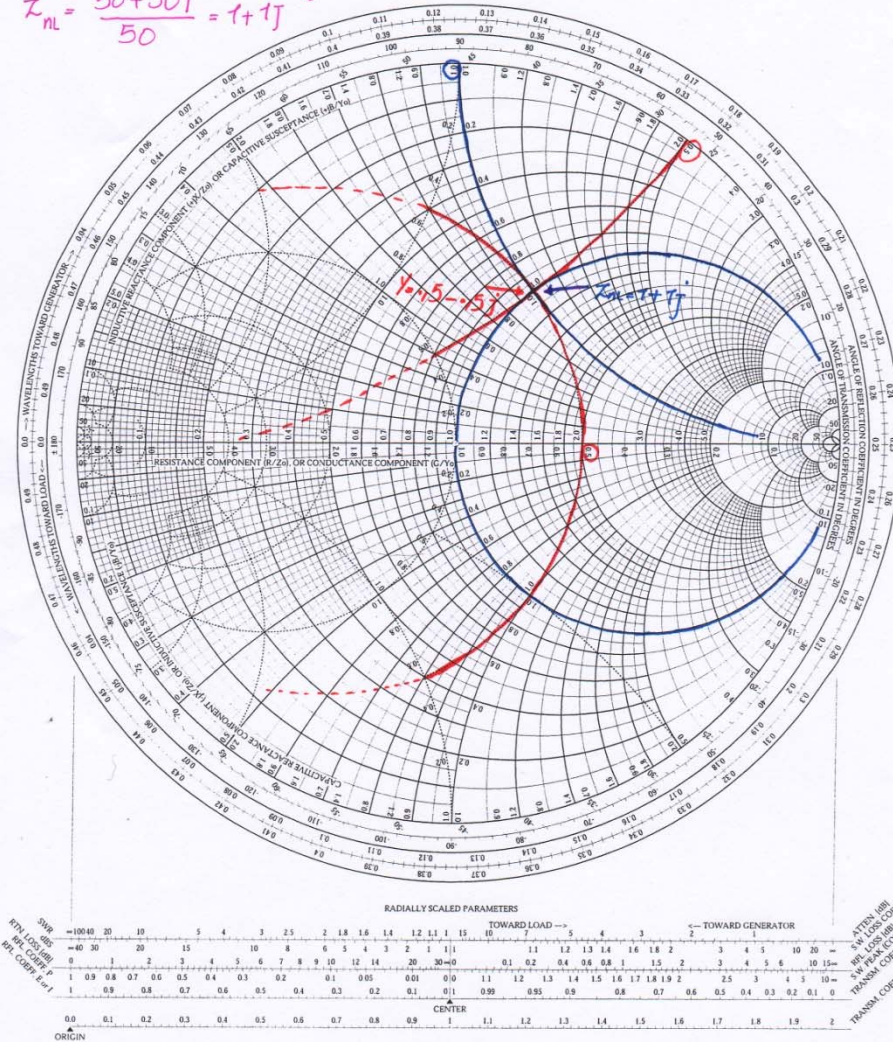


A ZY Smith Chart



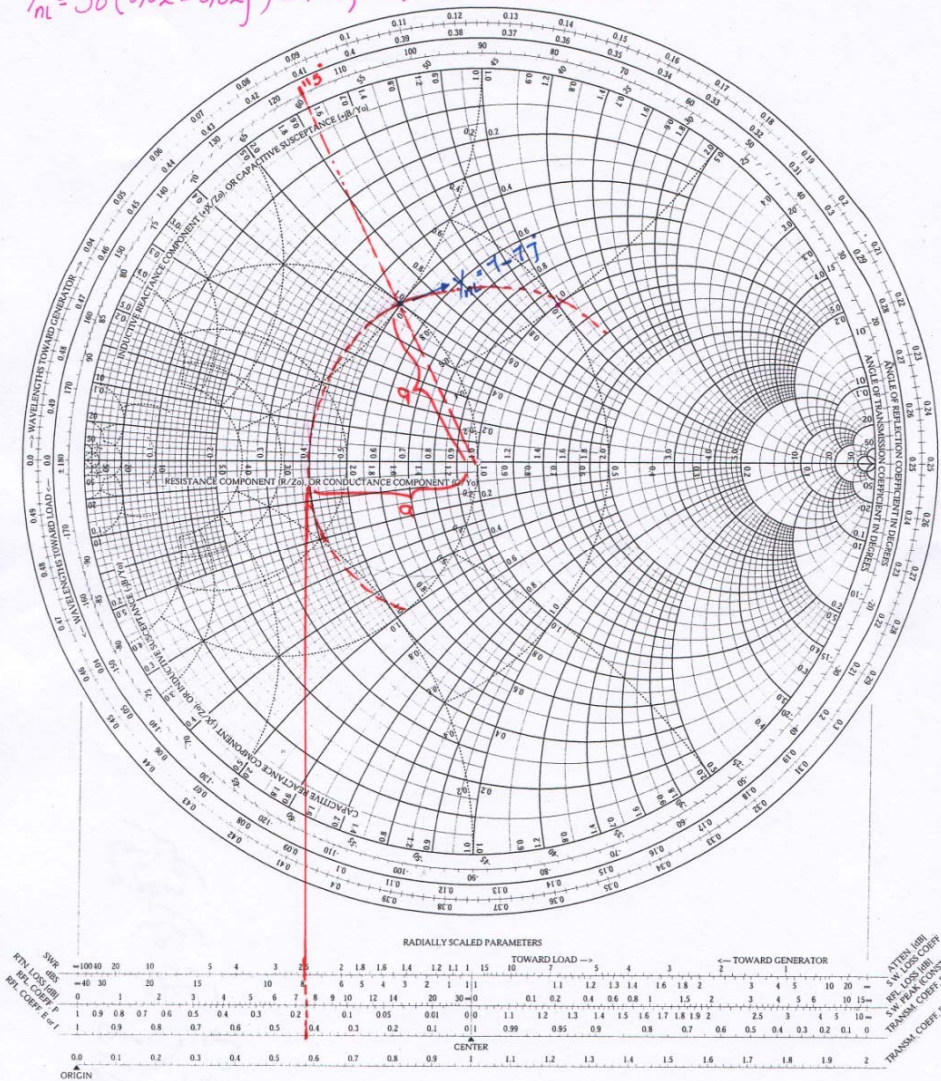
Smith Chart

ex) $Z = 50 + 50j$ \rightarrow $Y = ?$
 by Smith chart
 $Z_{nl} = \frac{50 + 50j}{50} = 1 + 1j$



ex) $Y = 0,02 - 0,02j \rightarrow \Gamma = ?$ Smith Chart

$Y_{in} = 50(0,02 - 0,02j) = 1 - 1j \Rightarrow \Gamma = 0,47 \angle 115^\circ$

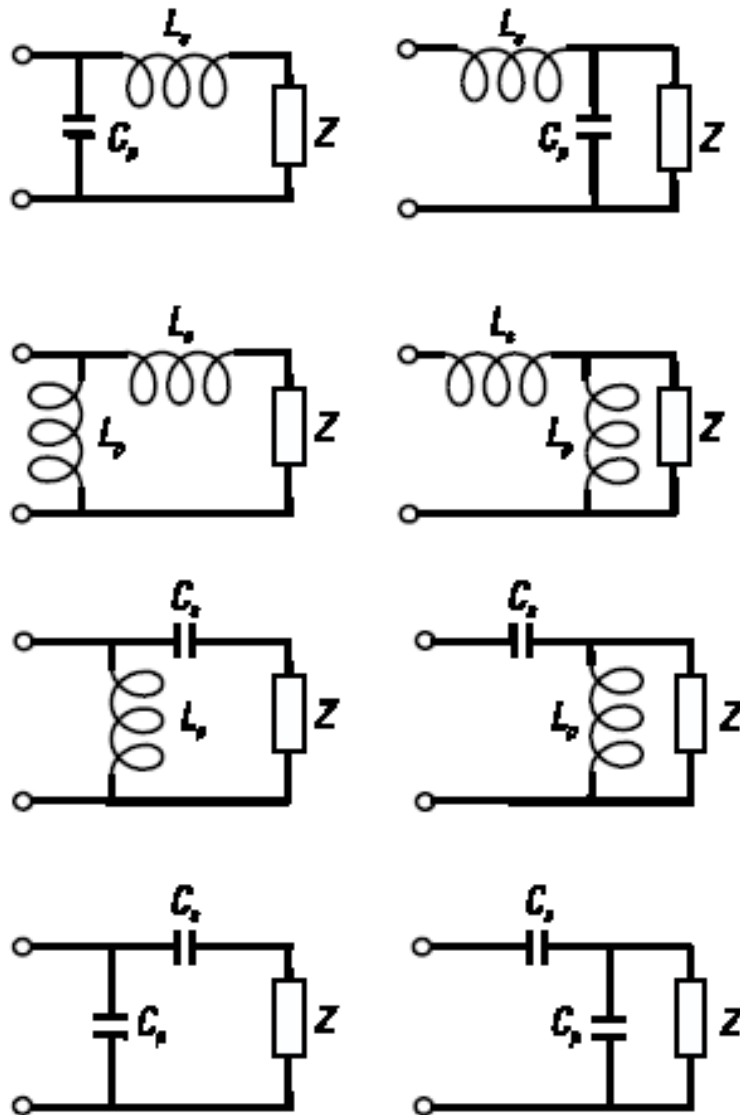


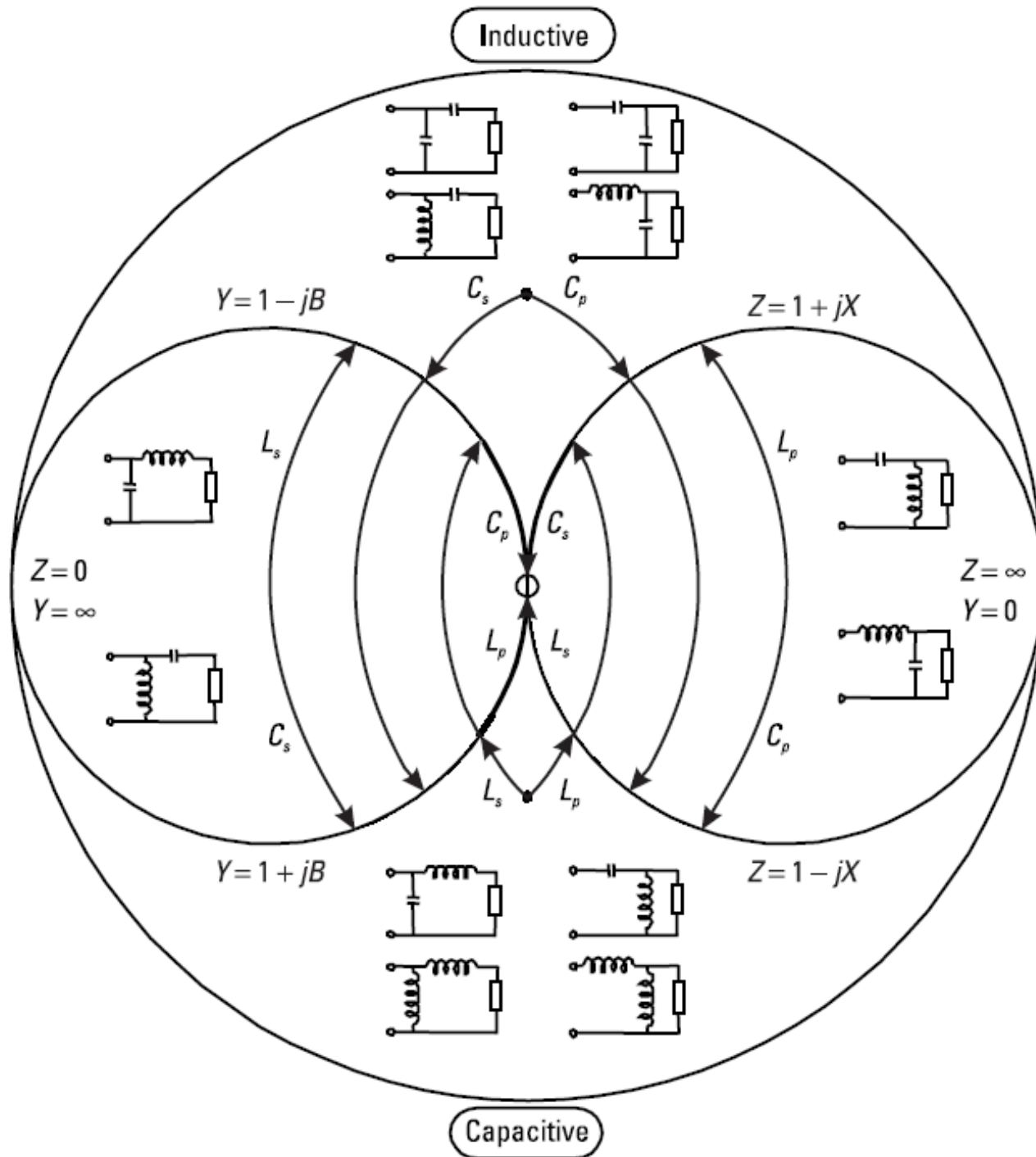
Impedance Matching

Using Lumped Components to Match Circuits

Component Added	Effect	Description of Effect
Series inductor	$z \rightarrow z + j\omega L$	Move clockwise along a resistance circle
Series capacitor	$z \rightarrow z - j/\omega C$	Smaller capacitance increases impedance ($-j/\omega C$) to move counterclockwise along a resistance circle
Parallel inductor	$y \rightarrow y - j/\omega L$	Smaller inductance increases admittance ($-j/\omega L$) to move counterclockwise along a conductance circle
Parallel capacitor	$y \rightarrow y + j\omega C$	Move clockwise along a conductance circle

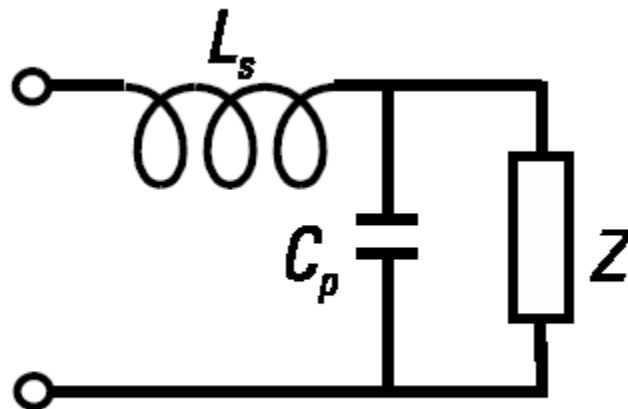
Eight Possible Impedance-Matching Networks





Example#1

Match $Z = 150 - 50j$ to 50Ω using the techniques just developed.



Normalized Impedance: $3 - 1j$

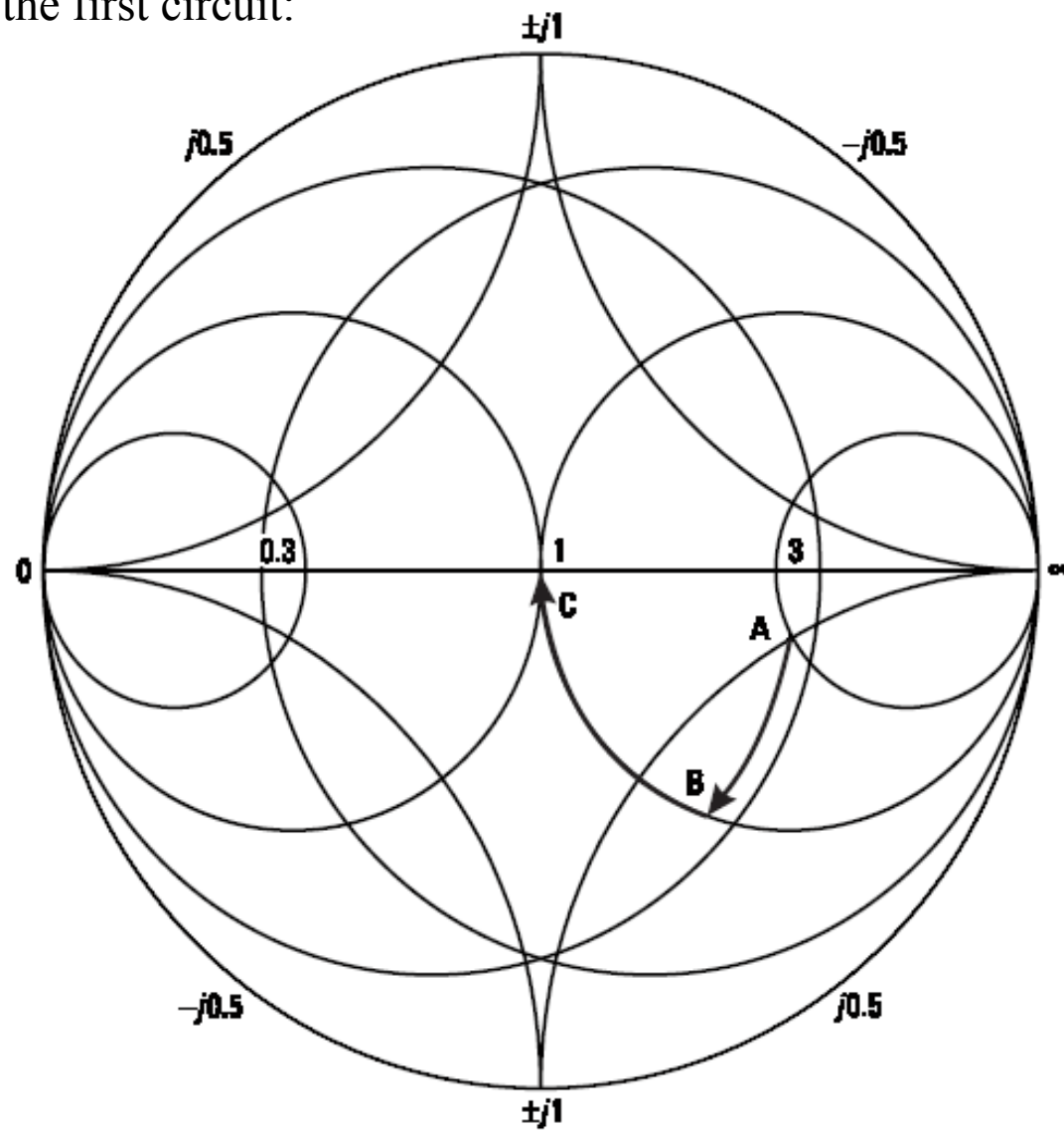
$$Y_A = 1/Z_A = 0.3 + j0.1$$

$Y_B = 0.3 + 0.458j$ we need a capacitor admittance of $0.348j$.

Since $Z_B = 1/Y_B = 1 - 1.528j$, an inductor reactance of $1.528j$ is needed to bring it to the center.

Cont'd

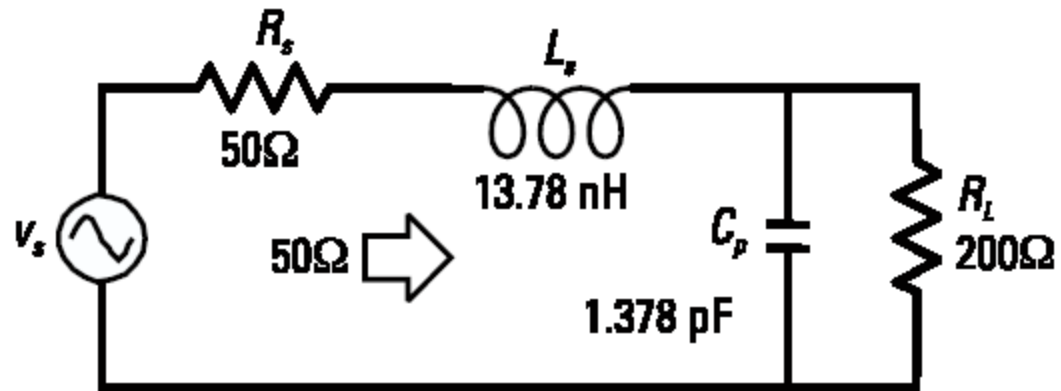
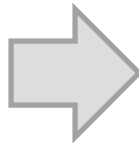
Match process of the first circuit:



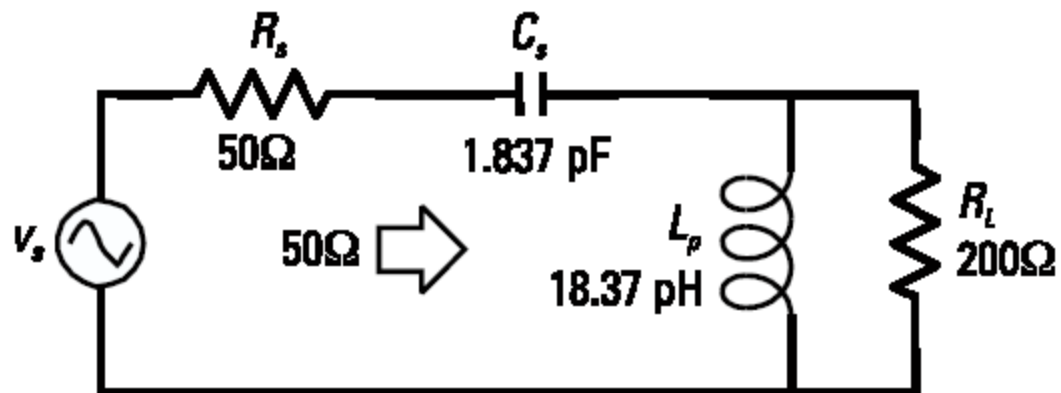
Example#2

Match a $200\text{-}\Omega$ load to a $50\text{-}\Omega$ source at 1 GHz

Low-pass structure

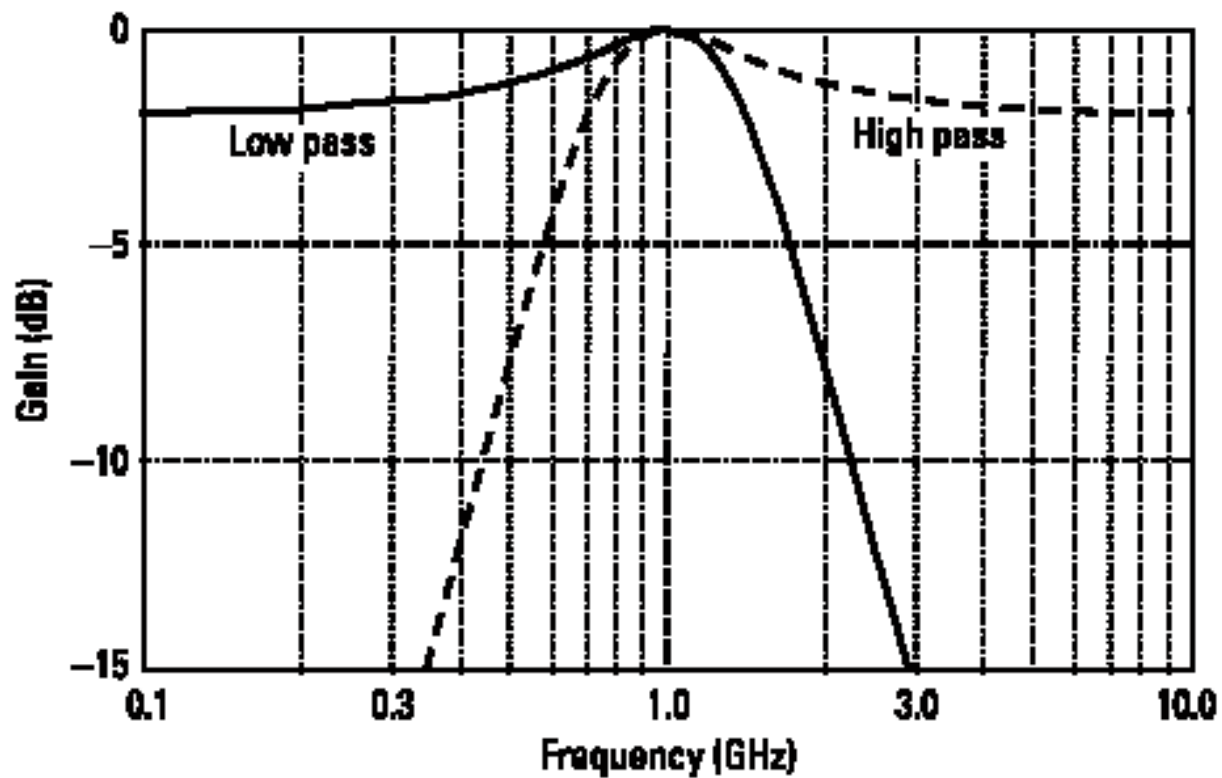


High-pass structure



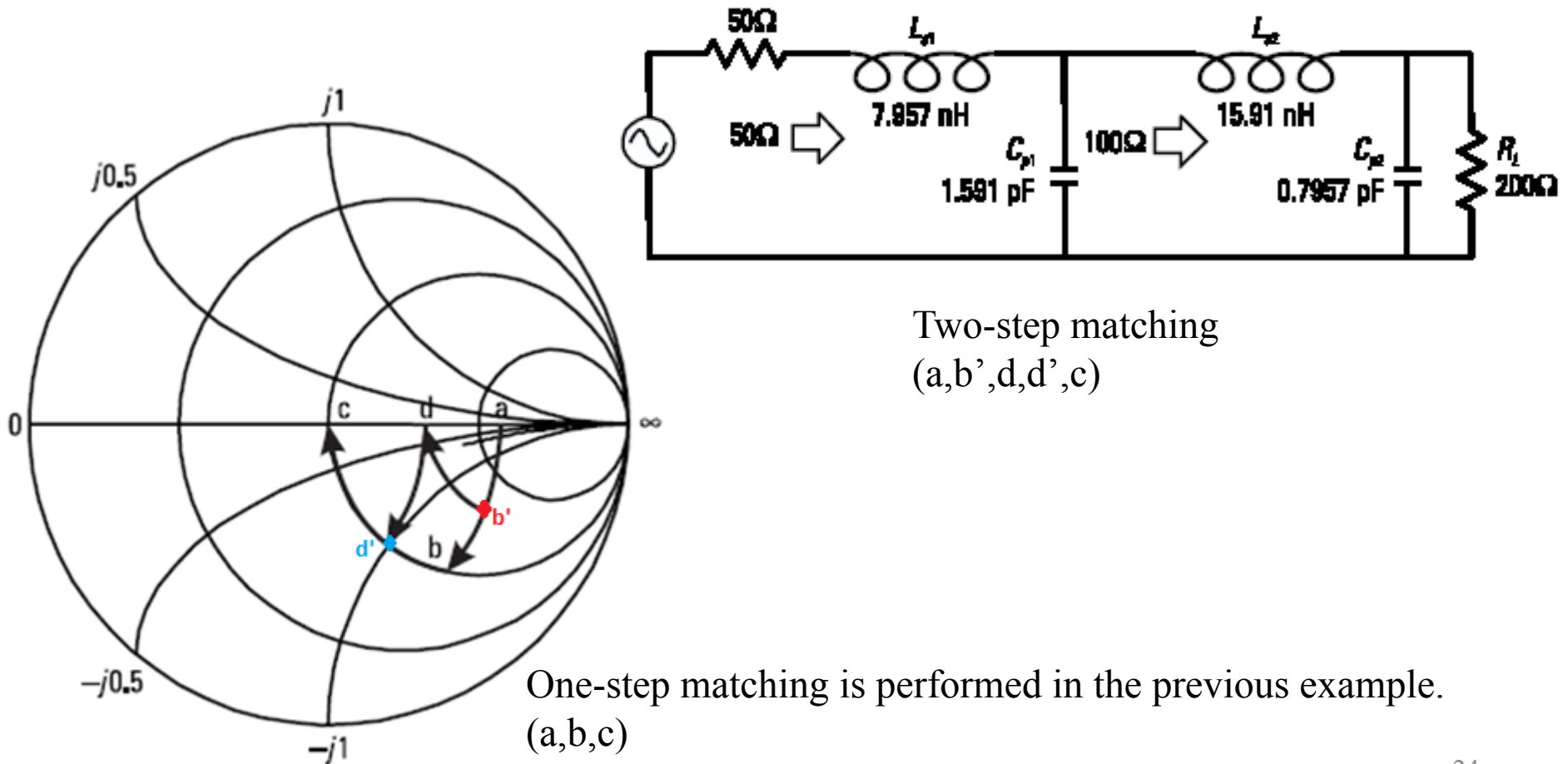
Cont'd

Frequency response for lowpass and highpass matching networks.



Example#3

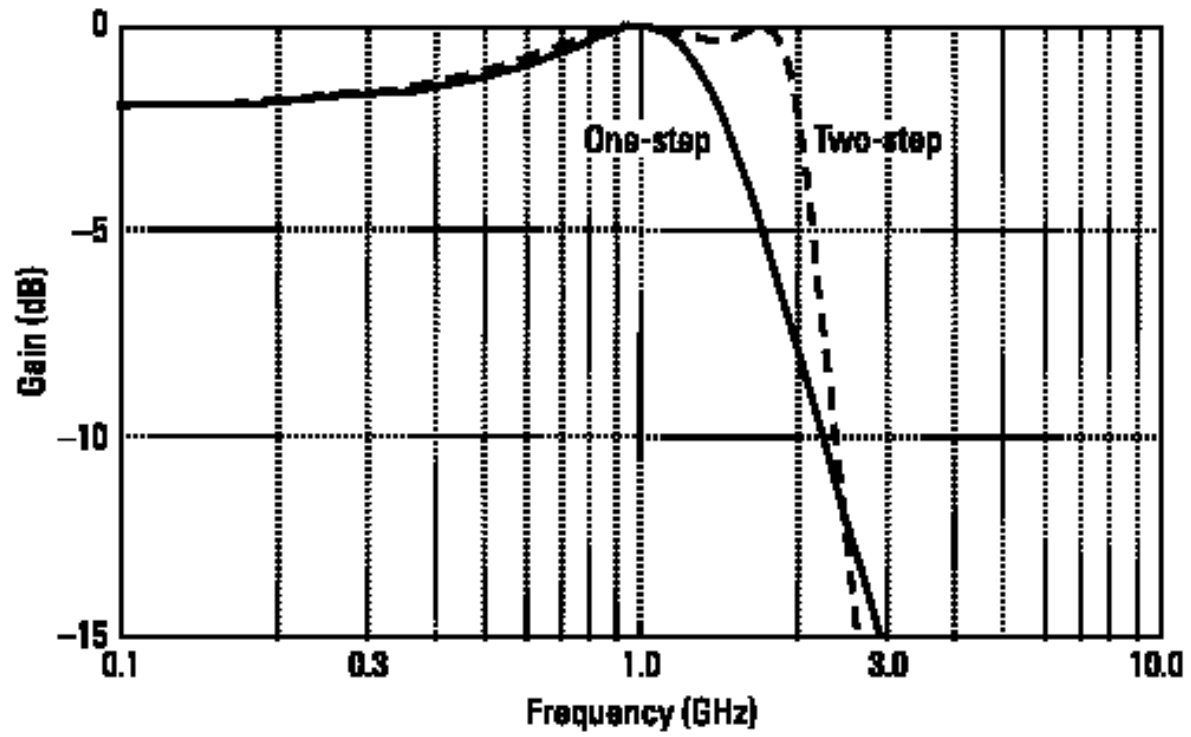
Match 200Ω to 50Ω at 1 GHz using an ell matching network. Do it first in one step, then do it in two steps matching it first to 100Ω . Compare the bandwidth of the two matching networks.



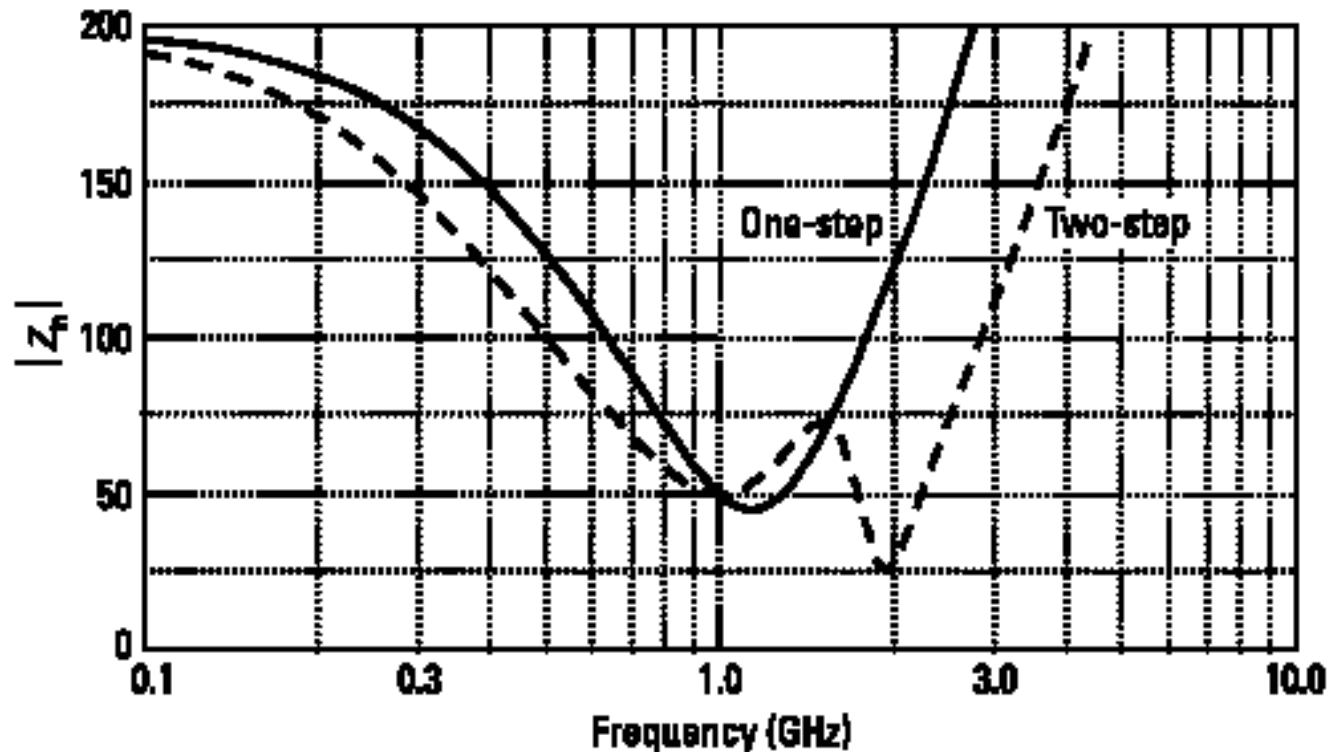
Two-step matching
(a,b',d,d',c)

One-step matching is performed in the previous example.
(a,b,c)

Cont'd

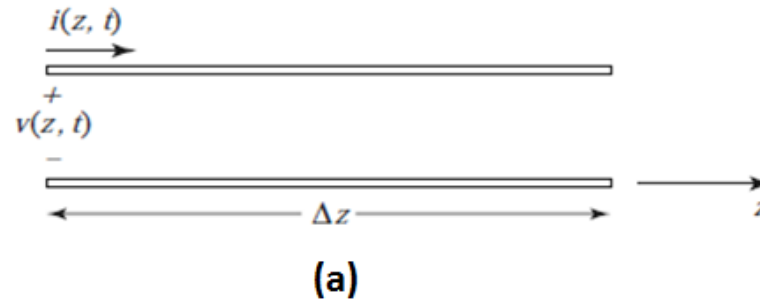


Cont'd



A comparison of frequency response shown in above figures clearly shows the bandwidth broadening effect of matching in two steps.

Circuit Model for Transmission Line

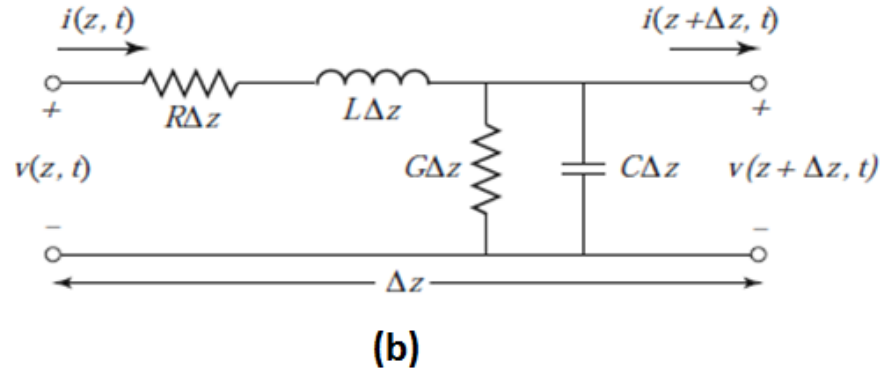


R: (Ω/m)

L: (H/m)

C: (F/m)

G: (S/m)



$$\text{KVL: } v(z,t) - R\Delta z i(z,t) - L\Delta z \frac{\partial i(z,t)}{\partial t} - v(z+\Delta z,t) = 0 \quad \Rightarrow \quad \frac{\partial v(z,t)}{\partial z} = -Ri(z,t) - L \frac{\partial i(z,t)}{\partial t}$$

$$\text{KCL: } i(z,t) - G\Delta z v(z+\Delta z,t) - C\Delta z \frac{\partial v(z+\Delta z,t)}{\partial t} - i(z+\Delta z,t) = 0 \quad \Rightarrow \quad \frac{\partial i(z,t)}{\partial z} = -Gv(z,t) - C \frac{\partial v(z,t)}{\partial t}$$

Sinusoidal steady-State Analysis:

$$\frac{dV(z)}{dz} = -(R+j\omega L)I(z)$$

$$\frac{dI(z)}{dz} = -(G+j\omega C)V(z)$$

Cont'd

$$\blacksquare \begin{cases} \frac{dV(z)}{dz} = -(R+j\omega L)I(z) \\ \frac{dI(z)}{dz} = -(G+j\omega C)V(z) \end{cases} \rightarrow \begin{cases} \frac{d^2V(z)}{dz^2} - \gamma^2V(z) = 0 \\ \frac{d^2I(z)}{dz^2} - \gamma^2I(z) = 0 \end{cases}$$

$$\gamma = \alpha + j\beta = \sqrt{(R + j\omega L)(G + j\omega C)}$$

$$\begin{aligned} V(z) &= V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} \\ I(z) &= I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z} \end{aligned}$$

or

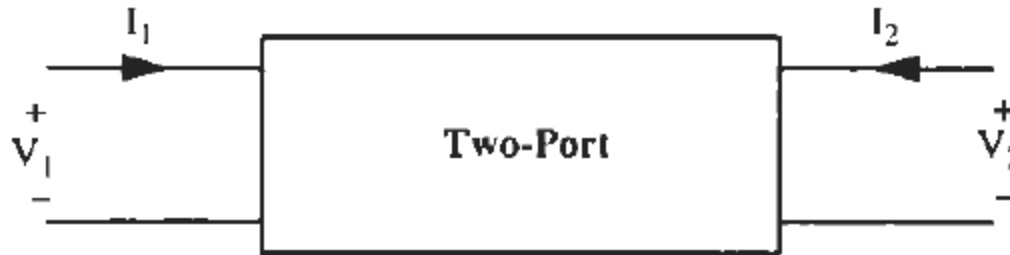
$$\begin{aligned} V(z) &= V_i e^{-\gamma z} + V_r e^{\gamma z} \\ I(z) &= I_i e^{-\gamma z} + I_r e^{\gamma z} \end{aligned}$$

It is easily proven that: $I(z) = \frac{\gamma}{R+j\omega L} (V_0^+ e^{-\gamma z} - V_0^- e^{\gamma z})$

$$\text{Defining: } Z_0 = \frac{R+j\omega L}{\gamma} = \sqrt{\frac{R+j\omega L}{G+j\omega C}} \rightarrow I(z) = \frac{V_0^+}{Z_0} e^{-\gamma z} - \frac{V_0^-}{Z_0} e^{\gamma z}$$

$$v(z,t) = |V_0^+| \cos(\omega t - \beta z + \phi^+) e^{-\alpha z} + |V_0^-| \cos(\omega t + \beta z + \phi^-) e^{\alpha z}$$

Z-Parameters

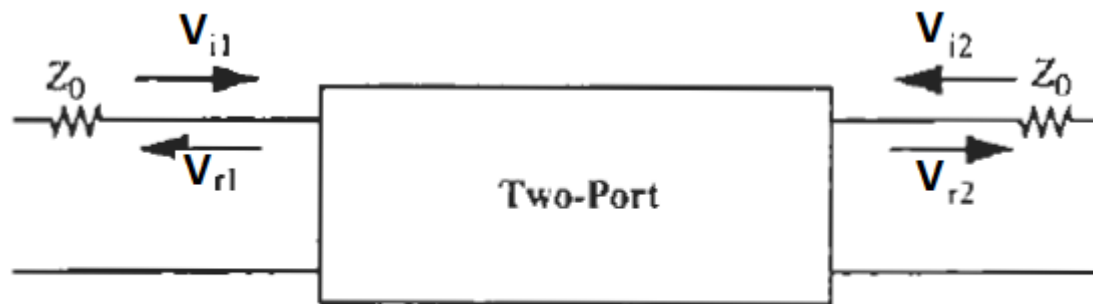


Port variable definitions.

$$V_1 = Z_{11}I_1 + Z_{12}I_2,$$

$$V_2 = Z_{21}I_1 + Z_{22}I_2.$$

S-Parameters



S-parameter port variable definitions.

$$V(z) = V_i e^{-\gamma z} + V_r e^{\gamma z}$$

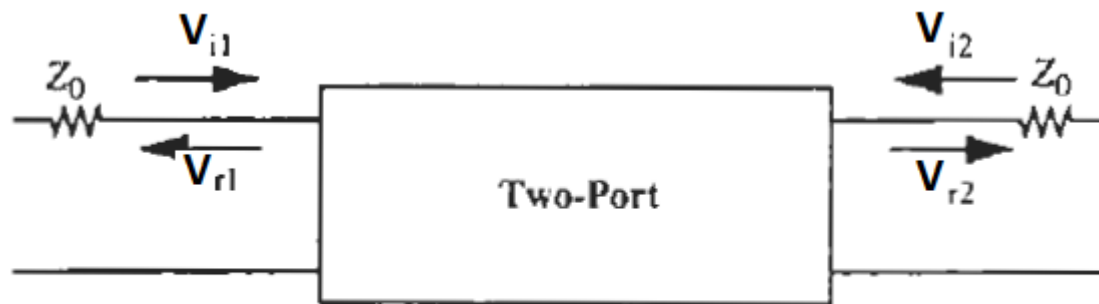
$$a_1 = \frac{V_{i1}}{\sqrt{Z_0}}, a_2 = \frac{V_{i2}}{\sqrt{Z_0}}, b_1 = \frac{V_{r1}}{\sqrt{Z_0}}, b_2 = \frac{V_{r2}}{\sqrt{Z_0}}$$

$$b_1 = s_{11}a_1 + s_{12}a_2$$

$$b_2 = s_{21}a_1 + s_{22}a_2$$

The normalization by the square root of Z_0 is a convenience that makes the square of the magnitude of the various a_n and b_n equal to the power of the corresponding incident or reflected wave.

S-Parameters



S-parameter port variable definitions.

Driving the input port with the output port terminated in Z_0 sets a_2 equal to zero, and allows us to determine the following parameters:

$$s_{11} = \frac{b_1}{a_1} = \frac{V_{r1}}{V_{i1}} = \Gamma_1,$$
$$s_{21} = \frac{b_2}{a_1} = \frac{V_{r2}}{V_{i1}}.$$

Thus, s_{11} is simply the input reflection coefficient, while s_{21} is a sort of gain since it relates an output wave to an input wave. Specifically, its magnitude squared is called the forward transducer power gain with Z_0 as source and load impedance.